An introduction to optimal transport (OT bootcamp)

Brendan Pass (U. Alberta)

June 20, 2022

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- Introduce some basic concepts from optimal transportation theory.
- Focus on ideas (rather than technical details) and on building intuition (with diagrams, non-rigorous proof sketches, etc.)
- Briefly cover a few topics requested by speakers.

Very incomplete list of references

- C. Villani. Topics in optimal transportation. AMS, 2003.
- C. Villani. Optimal transport: old and new. Springer, 2009.
- F. Santambrogio. Optimal transport for applied mathematicians. Birkhauser, 2015.
- G. Peyré and M. Cuturi. *Computational Optimal Transport: With Applications to Data Science* Now Publishers, 2019.
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- Given probability measures μ(x) (the source) and ν(y) (the target) on bounded domains X, Y ⊆ ℝⁿ, we say a map T : X → Y pushes μ forward to ν, and write T_#μ = ν, if μ(T⁻¹(A)) = ν(A) for all A ⊆ Y. We sometimes call these T's transport maps.
- Note: if dμ(x) = f(x)dx, dν(y) = g(y)dy, and T is a diffeomorphism (ie, 1 1, onto, smooth with a smooth inverse), this means T satisifies the change of variables equation f(x) = |detDT(X)|g(T(x)).

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- Given a cost function c(x, y), **Monge's optimal transport problem** is to minimize:

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- Example costs: $c(x, y) = |x y|, |x y|^2....$
- Challenging to analyze (lacks linearity, compactness...)

Leonid Kantorovich 1942: instead of sending all the mass at the source point x to target point y = T(x), allow **splitting**, so that the mass may be divided among **several** (or even infinitely many) target points.



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Intuitively, think of denoting the amount of mass moved from x to y by $\gamma(x, y)$.

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 $\gamma(X \times A) = \nu(A)$

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Marginals



 $\gamma(B \times Y) = \mu(B)$

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Marginals for a Monge type transport plan (transport map)



$$\nu(A) = \gamma(X \times A) = \gamma(B \times Y) = \mu(B) = \mu(T_{\text{cond}}^{-1}(A))$$

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 - Ex: For Monge type γ, the graph of T is the support (for a continuous T).

We say a set $S \subseteq X \times Y$ is *c*-monotone if for any (x_0, y_0) , $(x_1, y_1) \in S$ we have

$$c(x_0, y_0) + c(x_1, y_1) \le c(x_0, y_1) + c(x_1, y_0)$$

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• For
$$c(x, y) = |x - y|^2$$
, this amounts to $(x_1 - x_0) \cdot (y_1 - y_0) \ge 0$.

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Structure of solution: c-cyclical monotonicity

A set S ⊆ X × Y is c-cyclically monotone if for any finite collection of points (x₀, y₀), (x₁, y₁), ...(x_N, y_N) ∈ S we have

$$\sum_{i=0}^{N} c(x_i, y_i) \leq \sum_{i=0}^{N} c(x_i, y_{i+1}) + c(x_N, y_0)$$

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• For
$$c(x, y) = |x - y|^2$$
, this is

$$\sum_{i=0}^{N} x_i \cdot (y_i - y_{i+1}) + x_N \cdot (y_N - y_0) \ge 0$$

This is known simply as cyclical monotonicity.
• Suppose
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: $X, Y \subset \mathbb{R}$.

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 If (x₀, y₀), (x₁, y₁) are in the support of an optimal γ, c-monotonicity means that:

$$0 \geq c(x_1, y_1) + c(x_0, y_0) - c(x_0, y_1) - c(x_1, y_0)$$

=
$$\int_{y_0}^{y_1} \left[\frac{\partial c}{\partial y}(x_1, y) - \frac{\partial c}{\partial y}(x_0, y) \right] dy$$

=
$$\int_{y_0}^{y_1} \int_{x_0}^{x_1} \frac{\partial^2 c}{\partial x \partial y}(x, y) dx dy$$

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 $\implies (x_1 - x_0)(y_1 - y_0) \ge 0$. That is, γ is concentrated on a **monotone increasing** set.

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For $d\mu(x) = f(x)dx$, the only transport plan with monotone increasing support is concentrated on the graph of $T : X \to Y$ defined by:

$$\int_{-\infty}^{x} f(t)dt = \int_{-\infty}^{T(x)} g(s)ds$$



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- For probabilistically minded people, this is $T = (F_{\nu})^{-1} \circ F_{\mu}$, where F_{ν} and F_{μ} are the cumulative distribution functions.

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- For probabilistically minded people, this is $T = (F_{\nu})^{-1} \circ F_{\mu}$, where F_{ν} and F_{μ} are the cumulative distribution functions.
- Differentiating $\int_{-\infty}^{x} f(t)dt = \int_{-\infty}^{T(x)} g(s)ds$ with respect to x yields an ODE for T.

$$f(x) = T'(x)g(T(x))$$

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• Kantorovich's optimal transport problem between μ and ν is dual to the problem of maximizing:

$$\int_X u(x)d\mu(x) + \int_Y v(y)d\nu(y)$$

among functions u on X and v on Y such that $u(x) + v(y) \le c(x, y)$ for all $(x, y) \in X \times Y$

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• Kantorovich duality theorem:

$$\max_{u+v\leq c}\int_{X} u(x)d\mu(x) + \int_{Y} v(y)d\nu(y) = \min_{\gamma\in\Gamma(\mu,\nu)}\int_{X\times Y} c(x,y)d\gamma(x$$

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- The "≤" direction is easy to prove (just integrate both sides of the constraint u(x) + v(y) ≤ c(x, y) against γ).
- The duality is a key tool in analysis of OT problems.

• Minimax theory for

$$H(\gamma, u, v) = \int_{X \times Y} [c(x, y) - u(x) - v(y)] d\gamma(x, y)$$

+
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- For fixed γ the unconstrained supremum of

 (u, v) → H(γ, u, v) is ∫_{X×Y} c(x, y)dγ(x, y) if γ ∈ Γ(μ, ν),
 and ∞ otherwise.
 - So $\inf_{\gamma} \sup_{(u,v)} H(\gamma, \mu, \nu) = \inf_{\gamma \in \Gamma(\mu,\nu)} \int_{X \times Y} c(x, y) d\gamma(x, y)$ (the Kantorovich primal problem)

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- For fixed (u, v), the unconstrained infimum of γ → H(γ, u, v) is ∫_X u(x)dµ(x) + ∫_Y v(y)dν(y) if c(x, y) ≥ u(x) + v(y) everywhere and -∞ otherwise.

 So sup_(u,v) inf_Y H(γ, µ, ν) =

 $\sup_{u+v \leq c} \int_X u(x) d\mu(x) + \int_Y v(y) d\nu(y)$ (the dual problem)

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 - $c(x,y) \ge u(x) + v(y)$ everywhere and $-\infty$ otherwise.
 - So $\sup_{(u,v)} \inf_{\gamma} H(\gamma, \mu, \nu) =$ $\sup_{u+v \le c} \int_X u(x) d\mu(x) + \int_Y v(y) d\nu(y)$ (the dual problem)
- Applying a minimax theorem, $\inf_{\gamma} \sup_{(u,v)} H(\gamma, \mu, \nu) = \sup_{(u,v)} \inf_{\gamma} H(\gamma, \mu, \nu)$ gives duality.

More on duality: key facts

• Suppose that u(x) and v(y) solve the dual problem, and $\gamma(x, y)$ solves the primal. Then, since $u(x) + v(y) \le c(x, y)$ everywhere, but

$$\int_{X \times Y} [u(x) + v(y)] d\gamma(x, y) = \int_X u(x) d\mu(x) + \int_Y v(y) d\nu(y)$$
$$= \int_{X \times Y} c(x, y) d\gamma(x, y)$$

we must have

$$u(x)+v(y)-c(x,y)=0$$

 γ -almost everywhere.

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 γ -almost everywhere.

This means that x → u(x) - c(x, ȳ) is maximized at any x = x̄ such that (x̄, ȳ) ∈ spt(γ) (that is, any x̄ that gets transported to ȳ by γ). So

$$abla u(ar{x}) =
abla_{\mathsf{x}\mathsf{x}} c(ar{x},ar{y}) ext{ and } D^2 u(ar{x}) \leq D^2_{\mathsf{x}\mathsf{x}} c(ar{x},ar{y})$$

Brenier's Theorem: Suppose that μ is absolutely continuous with respect to Lebesgue measure and that $c(x, y) = \frac{|x-y|^2}{2}$. Then the solution γ to the Kantorovich problem is unique and concentrated on the graph of a function $T : X \to Y$. Furthermore, $T(x) = \nabla \phi(x)$ for a convex function ϕ , and T is the unique solution to the Monge problem.

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Remarks:

- In n = 1 dimension, convexity means that T'(x) = φ"(x) ≥ 0, so that T is monotone increasing – we recover our earlier result.
- A non-trivial part of the theorem is that there exists a (unique) convex function ϕ so that $\nabla \phi_{\#} \mu = \nu$. This fact by itself has many applications.

Let γ solve the primal problem and u(x), v(y) solve the dual.
 We have, for any (x, y) ∈ spt(γ),

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- The argument above shows that γ_0 and γ_1 concentrate on graphs of functions T_0 and T_1 . $\gamma_{1/2}$ then concentrates on the union of theses two graphs.

Graphical supports and their union



Graphical supports and their union

 $graph(v_0) \vee graph(v_1)^{-1}$ = $Spt(v_{11+})$

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- So we must have y = x - ∇u(x) = ∇(^{|x|²}/₂ - u)(x) := ∇φ(x) := T(x).
 Now, D²u(x) ≤ D²_{xx}c(x, y) = I, so D²φ(x) = D²(^{|x|²}/₂ - u)(x) = I - D²u(x) ≥ 0. So φ is convex.
- What about uniqueness? If γ_0 and γ_1 both solve Kantorovich's problem, so does $\gamma_{1/2} = \frac{1}{2}[\gamma_0 + \gamma_1]$, by linearity.
- The argument above shows that γ_0 and γ_1 concentrate on graphs of functions T_0 and T_1 . $\gamma_{1/2}$ then concentrates on the **union** of theses two graphs.
- But $\gamma_{1/2}$ concentrates on a graph, too (using the argument above again), which is impossible unless $T_0 = T_1$ in which case $\gamma_0 = \gamma_1$.

Remarks

The fact that the solution is graphical and unique doesn't really rely on *c* being quadratic. The same conclusions can be drawn for more general costs satisfying the **twist condition**, which is injectivity of *y* → ∇_x*c*(*x*, *y*) for each fixed *x*.

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$$c(x, y) = \frac{|x-y|^2}{2}$$
, note that $DT(x) = D^2\phi(x)$, so
 $\frac{f(x)}{g(\nabla\phi(x))} = |det(DT(x))| = detD^2\phi(x)$. That is, ϕ solves a Monge-Ampere equation.

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 Instead of using duality, one could use Rockafellar's theorem (a set S ⊆ ℝⁿ × ℝⁿ is cyclically monotone if any only if it is contained in the subdifferential of a convex function).

Brenier map: examples

• If μ is uniform on a ball, B(0,1) and ν uniform on the corresponding sphere, $\partial B(0,1)$, then

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$$\phi(x) = \overline{y} \cdot x + \frac{1}{2} x^T \Sigma^{1/2} x$$
, and $\nabla \phi(x) = \overline{y} + \Sigma^{1/2} x$

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Isoperimetric inequality: The surface area of any set $M \subseteq \mathbb{R}^n$ is greater than or equal to the surface area of a ball with the same volume.

$$Vol(M) = Vol(B_R(0)) \implies S(M) \ge S(B_R(0))$$

Isoperimetric inequality







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Proof:

• Take
$$f(x) = \chi_M, g(y) = \chi_{B_R(0)}$$
.

- $\nabla \phi(x)$ the Brenier map $\implies det(D^2 \phi(x)) = f(x)/g(\nabla \phi(x)) = 1$ (change of variables).
- Arithmetic mean dominates geometric mean (as ϕ is convex, $D^2\phi$ has positive eigenvalues) $\implies det^{1/n}(D^2\phi(x)) \leq \frac{1}{n}tr(D^2\phi(x)) = \frac{1}{n}\Delta\phi(x)$

$$\frac{1}{n}S(B_R(0))R = Vol(B_R(0)) = Vol(M)$$

$$= \int_M 1d^n x$$

$$= \int_M det^{1/n}(D^2\phi(x))dx$$

$$\leq \int_M \frac{1}{n}\Delta\phi(x)dx$$

$$= \frac{1}{n}\int_{\partial M} \nabla\phi(x) \cdot \vec{N}d^{n-1}S(x)$$

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Isoperimetric inequality



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Proof sketch:

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$$= \frac{1}{n}\int_{\partial M} \nabla\phi(x) \cdot \vec{N}d^{n-1}S(x)$$

$$\leq \frac{1}{n}\int_{\partial M} Rd^{n-1}S(x)$$

$$= \frac{1}{n}S(M)R$$

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Comments on the proof

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• The optimal transport proof is pretty **simple**; everything in the proof is first or second year mathematics (*except* Brenier's theorem)!

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- The optimal transport proof is pretty **simple**; everything in the proof is first or second year mathematics (*except* Brenier's theorem)!
- We prove an inequality about surfaces/curves/bodies in ℝⁿ by working with simple inequalities under the integral sign (geometric-arithmetic mean, Cauchy-Schwartz on ℝⁿ).
- This is a **common theme** in applications of optimal transport in geometry.

• Optimal transport can be used to derive a metric on the space of probability measures, which we call the Wasserstein distance:

$$W_2(\mu,
u) := \sqrt{\min_{\gamma \in \Gamma(\mu,
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- This is useful in a variety of applications when we want to compare two distributions of mass and the underlying distance plays a role.
- It works with discrete measures; in fact, x → δ_x isometrically embeds ℝⁿ into the space of probability measures.
- In one dimension, $W_2^2(\mu, \nu) = \int_0^1 |F_{\mu}^{-1}(t) F_{\nu}^{-1}(t)|^2 dt$, where F_{μ} and F_{ν} are the cdfs (ie, we compare μ and ν via their quantiles.)





Three probability measures: du(x) = f(x)dx, $d\nu(x) = g(x)dx$, $d\sigma(x) = h(x)dx$. Note that

$$||f - g||_{L^2} = ||f - h||_{L^2}$$

but

$$W_2(\mu, \nu) \ll W_2(\mu, \sigma)$$

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- Displacement interpolants are a natural way to interpolate (or average) between two probability measures, respecting the underlying geometry.
- The displacement interpolant between μ_0 and μ_1 is the curve of measures $\mu_t := ((1 - t)I + t\nabla\phi)_{\#}\mu_0$, where $\nabla\phi$ is the Brenier (optimal transport) map between μ_0 and μ_1 .



Probability measures μ_0 and μ_1 .



Probability measures μ_0 and μ_1 . The linear interpolant $\mu_t = (1-t)\mu_0 + t\mu_1$



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 In general, displacement interpolation tends to preserve shapes/geometric features better than other interpolation methods.

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- For an absolutely continuous probability measure at time t, $d\mu_t(x) = f_t(x)dx$ with density $f_t(x)$, we let $v_t(x)$ be the velocity of a particle at point x and time t.

- There is also a way to view the Wasserstein distance through action minimizing curves in the space of probability measures.
- For an absolutely continuous probability measure at time t, $d\mu_t(x) = f_t(x)dx$ with density $f_t(x)$, we let $v_t(x)$ be the velocity of a particle at point x and time t.
 - The divergence ∇ · (v_t(x)f_t(x)) tells us how much mass is moving away from, or towards, the point x at time t.
 - Conservation of mass gives us the continuity equation: $f'_t(x) + \nabla \cdot (v_t(x)f_t(x)) = 0.$

- Recall that in ℝⁿ, the squared distance |x₀ x₁|² is given by min ∫₀¹ |v_t|²dt, where the minimum is among curves x_t joining x₀ and x₁ with velocity v_t = x'_t.
 - The optimal curve is the straight line x_t = x₀ + t(x₁ x₀) (so that v_t = x₁ x₀ is constant).

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- Analagously, the squared Wasserstein distance between probability measures μ_0 and μ_1 may be written

$$W_2^2(\mu_0,\mu_1) = \min \int_0^1 \int_{\mathbb{R}^n} f_t(x) |v_t(x)|^2 dx dt$$

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• The optimal velocity, $v_t(x) = \nabla \phi(x) - x$ is exactly the Brenier map.

Wasserstein gradient flows

• Certain evolution PDEs can be interpreted as gradient flows in the Wasserstein space: that is, a distribution of mass is rearranging itself so as to decrease a certain functional as quickly as possible, relative to the Wasserstein metric.
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- Recall the gradient flow on \mathbb{R}^n for a function $F : \mathbb{R}^n \to \mathbb{R}$ is

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• How can this idea be adapted to Wasserstein space?

• First, we have to understand gradients. The key point from \mathbb{R}^n we want to translate is that for a curves x_t , $\frac{d}{dt}(F(x(t))) = \nabla F(x(t)) \cdot x'(t)$

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- For a curve of measures, $d\mu_t(x) = f_t(x)dx$, the classical way (working in L^2) to interpet its velocity would be as $f'_t(x)$ (the rate of change of mass at the point x).

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- Optimal transport is about how mass *moves*; we therefore interpret the velocity of a curve by its velocity vector at individual points, given by f'_t(x) = −∇ · (v_t(x)f_t(x))

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- Optimal transport is about how mass *moves*; we therefore interpret the velocity of a curve by its velocity vector at individual points, given by f'_t(x) = −∇ · (v_t(x)f_t(x))
- Therefore, for a functional \mathcal{F} on the space of probability measures, we would like ot define $\nabla \mathcal{F}$ so that for each curve $d\mu_t(x) = f_t(x)dx$, we have

$$\frac{d}{dt}(\mathcal{F}(\mu_t)) = \langle w_t, v_t \rangle_{L^2(\mu_t)} = \int_{\mathbb{R}^n} w_t(x) \cdot v_t(x) d\mu_t(x)$$

where $f_t' = -\nabla \cdot (v_t f_t)$ and $(\nabla \mathcal{F})(\mu_t) = -\nabla \cdot (w_t f_t)$

 \bullet For simplicity, consider functionals ${\cal F}$ on Wasserstein space of the form

$$\mathcal{F}(\mu) = \int_{\mathbb{R}^n} U(f(x)) dx$$

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• We would like to write this as an integral against the corresponding velocity vector *v*_t which satisfies:

$$f'_t(x) = -\nabla \cdot (v_t(x)f_t(x))$$

$$\begin{aligned} \frac{d}{dt}\mathcal{F}(\mu_t) &= \int_{\mathbb{R}^n} U'(f_t(x))f'_t(x)dx \\ &= -\int_{\mathbb{R}^n} U'(f_t(x))\nabla \cdot (v_t(x)f_t(x))dx \\ &= \int_{\mathbb{R}^n} \nabla (U'(f_t(x))) \cdot v_t(x)f_t(x)dx \end{aligned}$$

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This establishes $\nabla(U'(f_t(x)))$ as the **velocity** corresponding to the gradient. The corresponding rate of change of the density then comes from the continuity equation:

$$(\nabla \mathcal{F})(\mu_t) = -\nabla \cdot (f_t(x)\nabla (U'(f_t(x))).$$

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$$(\nabla \mathcal{F})(\mu_t) = -\nabla \cdot (f_t(x)\nabla (U'(f_t(x))).$$

So, the gradient of \mathcal{F} evaluated at $d\mu(x) = f(x)dx$ is $-\nabla \cdot (f(x)\nabla(U'(f(x))))$.

Take the entropy: $\mathcal{F} = \int_{\mathbb{R}^n} f(x) \log(f(x)) dx$ (so $U(r) = r \log(r)$). Then its Wasserstein gradient is

$$\begin{aligned} -\nabla \cdot (f(x)\nabla(U'(f(x))) &= -\nabla \cdot (f(x)\nabla(\log(f(x)+1))) \\ &= -\nabla \cdot (f(x)(\frac{\nabla f(x)}{f(x)})) = -\Delta f(x) \end{aligned}$$

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So the Wasserstein gradient flow of the entropy is given by:

$$f'(x) = \Delta f(x)$$

The heat equation!

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The **heat equation!** (Note: it is actually easy to show that the entropy decreases for solutions of the heat equation – this show that it decreases as quickly as possible relative to the Wasserstein metric. On the other hand, it is well known that the heat equation is the gradient flow of the Dirichlet energy, $\int_X |\nabla f(x)|^2$ under the L^2 metric.)

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The **heat equation!** (Note: it is actually easy to show that the entropy decreases for solutions of the heat equation – this show that it decreases as quickly as possible relative to the Wasserstein metric. On the other hand, it is well known that the heat equation is the gradient flow of the Dirichlet energy, $\int_X |\nabla f(x)|^2$ under the L^2 metric.) Many other examples (see, for example, Villani_TOT p.252).

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• For $x'(t) = -\nabla F(x(t))$ in \mathbb{R}^n , think of discretizing time $t_0, t_1, ...,$ with $h = t_{i+1} - t_i$.

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- For $x'(t) = -\nabla F(x(t))$ in \mathbb{R}^n , think of discretizing time $t_0, t_1, ...,$ with $h = t_{i+1} t_i$.
- Then $x(t_{i+1}) x(t_i) \approx h x'(t_{i+1}) = -h \nabla F(x_{t_{i+1}}).$

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- Similarly, on Wasserstein space, we expect choosing $\mu_{t_{i+1}}$ to minimize

$$\mu \mapsto \mathcal{F}(\mu) + \frac{1}{2h} W_2^2(\mu_{t_i}, \mu)$$

to be close to Wasserstein gradient flow (or the corresponding PDE).

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